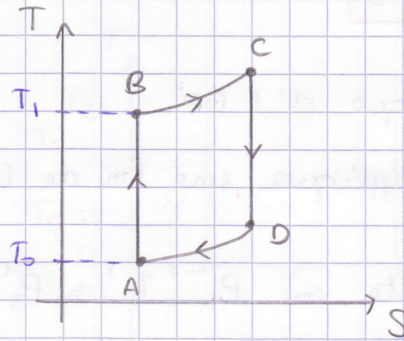
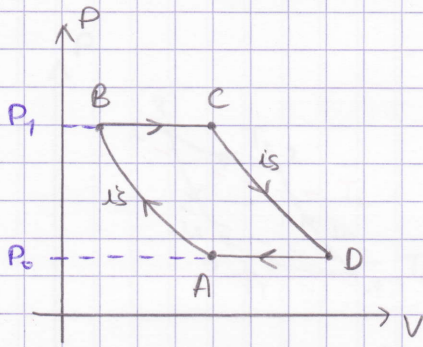


Ex TS-6. Turbomoteur.

1)



$$\delta = \{m, \text{GP}, \delta, \eta\}$$

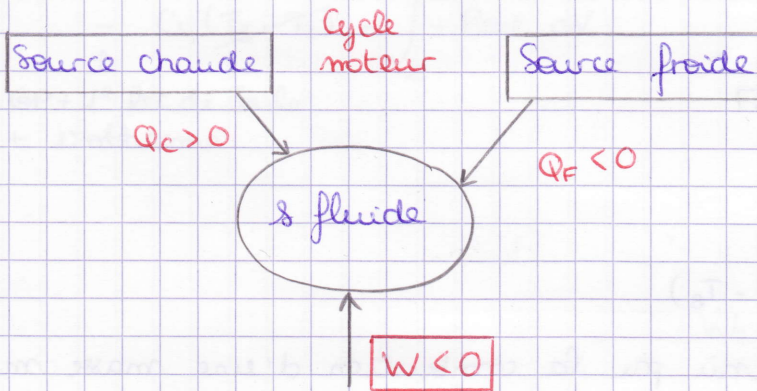
$$b = \frac{P_1}{P_0} = 5$$

Justification \rightarrow of laws.

2) Cycle moteur.

$$\eta = \frac{\text{gdein utile}}{\text{gdein coiteuse}} = \frac{-W}{Q_c} = 1 + \frac{Q_f}{Q_c}$$

$$(1P) W + Q_f + Q_c = 0$$



$$Q_c = Q_{AB} + Q_{BC}$$

$$Q_f = Q_{CD} + Q_{DA}$$

$$\text{car } \begin{cases} Q_{AB} = 0 \text{ (adiabatique "au" entropique)} \\ Q_{BC} = Q_{BC,p} = \Delta H_{BC} = C_p(T_c - T_b) > 0 \\ Q_{CD} = 0 \text{ (ad)} \\ Q_{DA} = Q_{DA,p} = \Delta H_{DA} = C_p(T_a - T_0) \end{cases}$$

\uparrow
2^e loi deoule pr 1 GP
 \downarrow

$$\eta = 1 + \frac{T_A - T_0}{T_C - T_B}$$

A → B : isentropique d'1 GP

↳ on peut appliquer une loi de Laplace.

$$P^{1-\gamma} T^\gamma = \text{cte} \Leftrightarrow P_A^{1-\gamma} T_A^\gamma = P_B^{1-\gamma} T_B^\gamma$$

$$T_A^\gamma = T_B^\gamma \left(\frac{P_B}{P_A} \right)^{\frac{1-\gamma}{\gamma}}$$

$$T_A = T_B (b)^{\frac{1-\gamma}{\gamma}} \quad \text{car} \quad \frac{P_B}{P_A} = \frac{P_1}{P_0} \equiv b$$

de m̂ par C → D : $P_0^{1-\gamma} T_0^\gamma = P_c^{1-\gamma} T_c^\gamma$

$$T_0 = T_c \left(\frac{P_0}{P_c} \right)^{\frac{1-\gamma}{\gamma}} = T_c (b)^{\frac{1-\gamma}{\gamma}} \quad \text{car} \quad \frac{P_0}{P_c} = \frac{P_1}{P_0} \equiv b$$

$$\eta = 1 + \frac{(T_0 - T_c) b^{\frac{1-\gamma}{\gamma}}}{T_c - T_B}$$

$$\eta = 1 - b^{\frac{1-\gamma}{\gamma}} = 0,37$$

3) $Q_c = \int Q_{c,p} = c_p (T_c - T_B)$

chaleur fournie par la combustion d'une masse m_1 de carburant

$$Q_c = \begin{cases} \frac{\gamma m R}{\gamma - 1} (T_c - T_B) = \frac{m}{\eta} \frac{\gamma R}{\gamma - 1} (T_c - T_B) = m c_p (T_c - T_B) \\ m_1 q \end{cases}$$

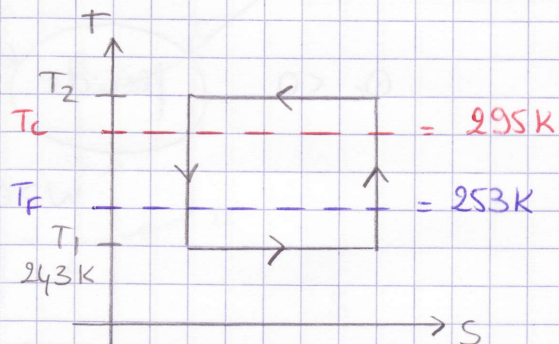
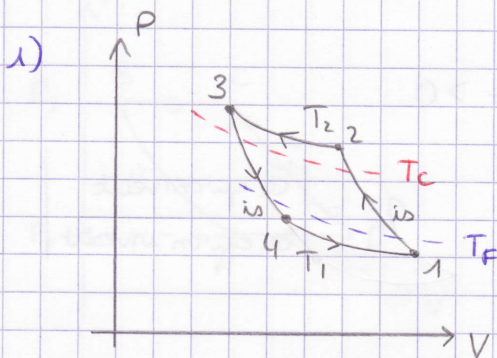
puissance calorifique massique du carburant

$$c_p = \frac{1}{\eta} \frac{\gamma R}{\gamma - 1} = 1,003 \text{ kJ} \cdot \text{K}^{-1} \cdot \text{kg}^{-1}$$

$$\rightarrow T_c = \frac{m_1 q}{m c_p} + T_B = \frac{m_1}{m} q \eta \frac{\gamma - 1}{\gamma R} + T_A b^{\frac{\gamma - 1}{\gamma}} \quad T_c = 2950 \text{ K}$$

avec $T_B = T_A b^{\frac{\gamma - 1}{\gamma}}$

Ex TS.8. Fugopompe.



Rq: Ce n'est pas le cycle de Carnot car les isothermes ne sont pas réversibles (si elles l'étaient il y aurait équilibre thermique avec les 2 sources à $T_1 = T_f$ et $T_2 = T_c$).

2) 2 adiabatiques : $Q_{12} = Q_{34} = 0$

$$Q_{23} = \Delta U_{23} - W_{23}$$

↑
(1P)

$$= C_v (T_3 - T_2) - \int_2^3 -P_{ext} dV$$

↑
GP + 1^{re} loi de Joule
+ isotherme

$$W_{23} = \int_2^3 \delta W = \int_2^3 -P_{ext} dV$$

$$\overset{TPS}{\curvearrowright} = \int_2^3 -P dV$$

$$\overset{GP}{\curvearrowright} = \int -\frac{nRT}{V}$$

$$\overset{isoth}{\curvearrowright} = -nRT_2 \int_2^3 \frac{dV}{V}$$

$$W_{23} = -nRT_2 \ln \frac{V_3}{V_2}$$

$$\text{or } \frac{V_3}{V_2} = \frac{nRT_3}{P_3} \frac{P_2}{nRT_2} = \frac{P_2}{P_3} = \frac{1}{3}$$

donc $W_{23} = nRT_2 \ln 3 > 0$ (compression isotherme)

$$\hookrightarrow Q_c = Q_{23} = -W_{23} = -nRT_2 \ln 3 < 0$$

